

# **Magnetic properties and electronic structures of rare-earth-transition metal compounds**

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1. Exchange enhanced paramagnets

2. Magnetic ordered compounds

2.1 Exchange interactions

2.2 Induced cobalt moments

2.3 Magnetism and pressure effects

Transition metals:

2.3.1 Strong ferromagnetism,  $\text{RCO}_5$  compounds

2.3.2 Weak ferromagnetism,  $\text{RCO}_2$  compounds

2.3.3 Intermediate degree of localization

$\text{RCO}_4\text{X}$  ( $\text{X} = \text{B, Si}$ ), R-Co-B

3. Conclusions

Methods

Crystal structures

Magnetic properties

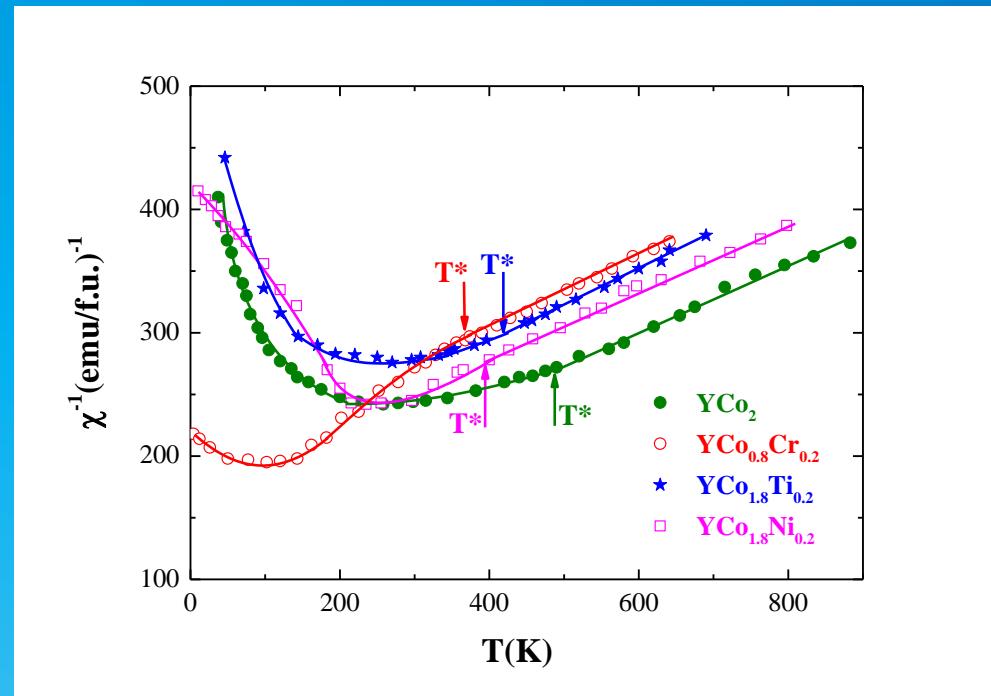
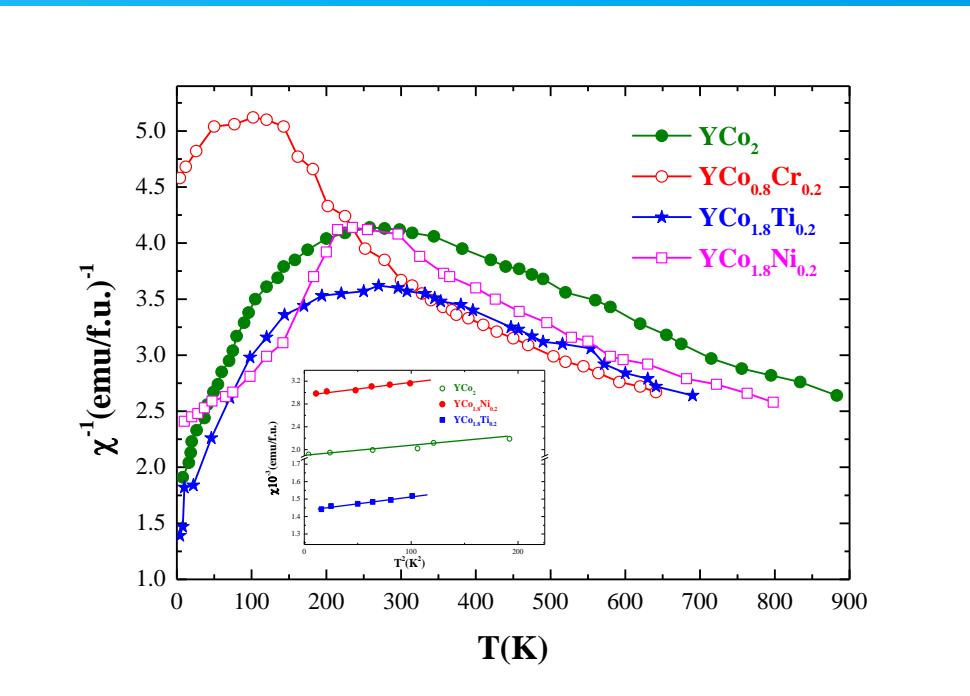
Magnetocaloric effect

Band structure calculations

LMTO-ASA

LDA+U

# 1. Exchange enhanced paramagnets



Low temperature  $\chi = \chi_0(1+aT^2)$

High temperature  $T > T^*$   $\chi = C(T-\theta)^{-1}$   
 $\theta < 0$

## Self consistent theory of spin fluctuations

Wave number dependent susceptibility,  $\chi_q$ , for a nearly ferromagnetic alloy has a large enhancement for small  $q$  values

$$\chi_q = \frac{\bar{\chi}_q}{1 - J\bar{\chi}_q(\mu_0 \mu_B^2)}$$

Frequency of longitudinal spin fluctuations  $\omega^* \propto \frac{1}{\tau}$   
τ-lifetime of LSF

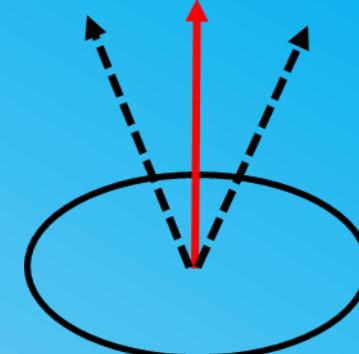
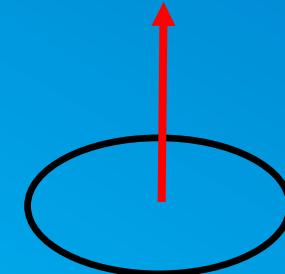
Low temperature

$$\omega^* > \frac{k_B T}{\hbar} \quad \text{thermal fluctuations (transversal) slow}$$

$$\chi = s\chi_p \left[ 1 + \frac{\pi^2}{6} \left( 2 \frac{\eta''}{\eta} - 1.2 \frac{\eta'^2}{\eta^2} \right)_{E_F} s^2 T^2 \right]$$

Approximation for nonmagnetic state

$$\chi \propto T^2 \quad \chi(T) \nearrow \quad \text{as } T \nearrow \\ \eta'' > 0 \text{ (necessary condition, not sufficient)}$$

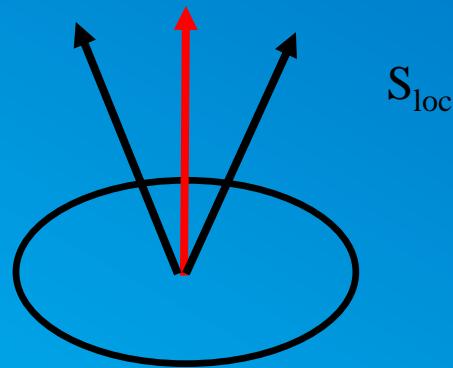


# High temperature

Average mean amplitude of LSF is temperature dependent

$$\left\langle S_{loc}^2 \right\rangle = 3k_B T \sum_q \chi_q$$

$$\omega^* < \frac{k_B T}{\hbar}$$

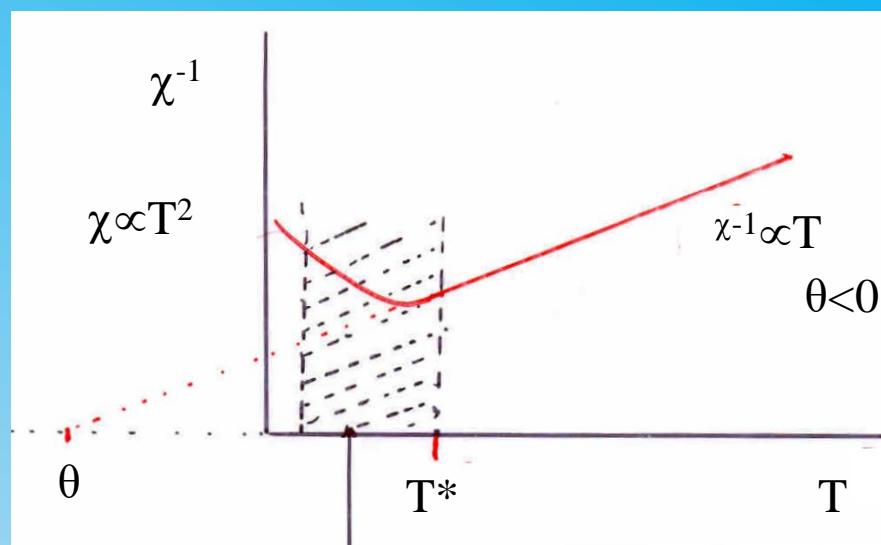


$S_{loc} \nearrow$  as  $T \nearrow$  up to  $T^*$  ( $S_{loc}$ )

$S_{loc}$  determined by charge neutrality condition

The system behaves as having local moments for temperatures  $T > T^*$  where the frequency of thermal fluctuations is higher than of longitudinal.

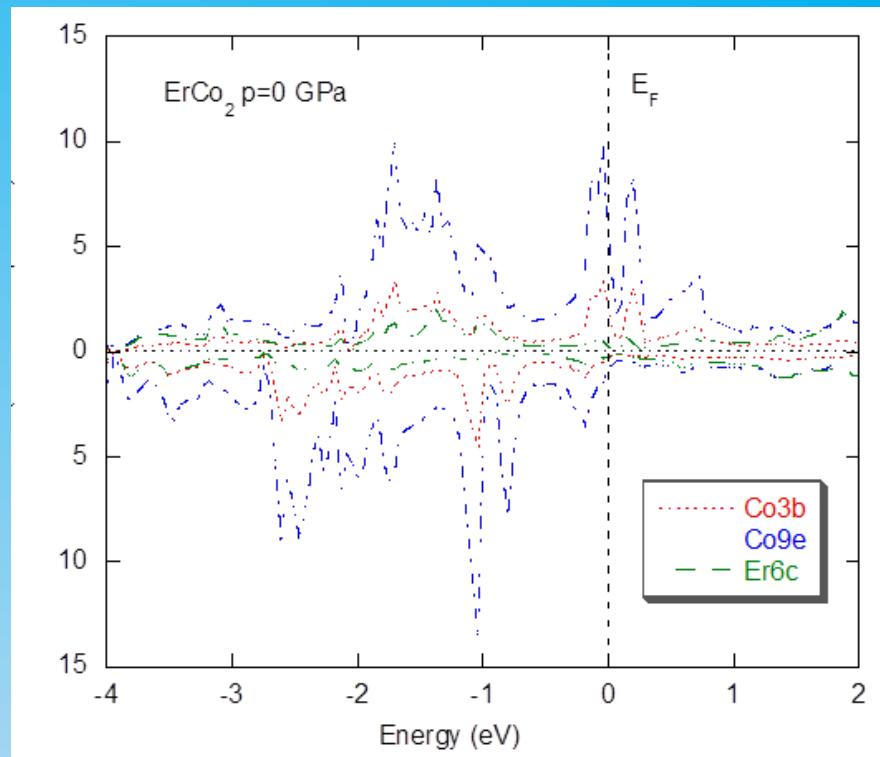
Transition from exchange enhanced paramagnetism to Curie-Weiss type behaviour.



Compound	$\chi_{\text{exp}} \cdot 10^{-3}$ at 2 K (emu/f.u.)	$\chi_{\text{calc}} \cdot 10^{-3}$ at 2 K (emu/f.u.)	$T_{\max}$ (K)	$T^*$ (K)	$a \cdot 10^{-3}$		$M_{\text{eff}} \text{Co}$ $\mu_B/\text{atom}$
	exp.	theor.					
LuCo <sub>2</sub>	1.82	1.92	370	550	0.764	0.91	4.10
YCo <sub>2</sub>	1.95	2.25	260	485	1.64	1.81	3.86
YCo <sub>1.8</sub> Ni <sub>0.2</sub>	2.9	3.02	215	408	1.24	1.36	3.84
YCo <sub>0.9</sub> Ti <sub>0.1</sub>	1.271		275	450	1.068		3.95
YCo <sub>1.875</sub> T <sub>0.125</sub>		1.796				0.9961	
YCo <sub>1.8</sub> Ti <sub>0.2</sub>	1.442		250	420	0.908		3.90
YC <sub>1.75</sub> Ti <sub>0.25</sub>		2.046				0.8945	
YCo <sub>0.8</sub> Cr <sub>0.2</sub>	4.58		180	370			

## 2. Magnetic ordered compounds

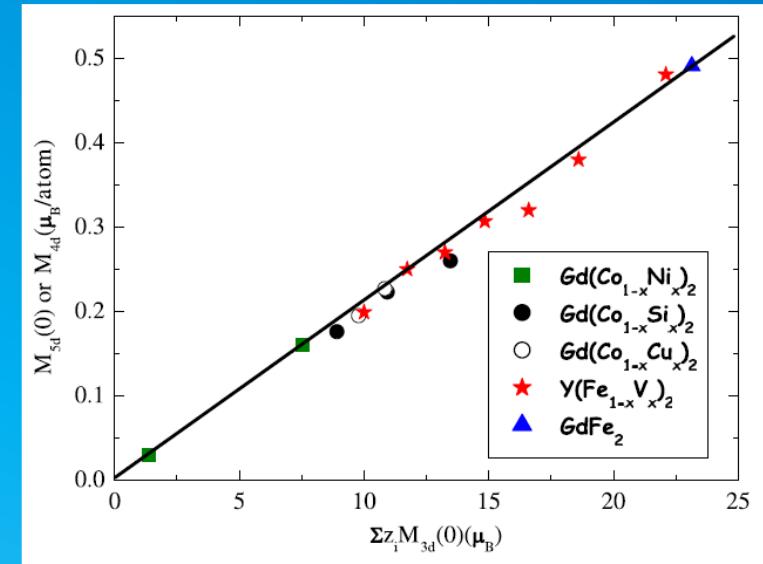
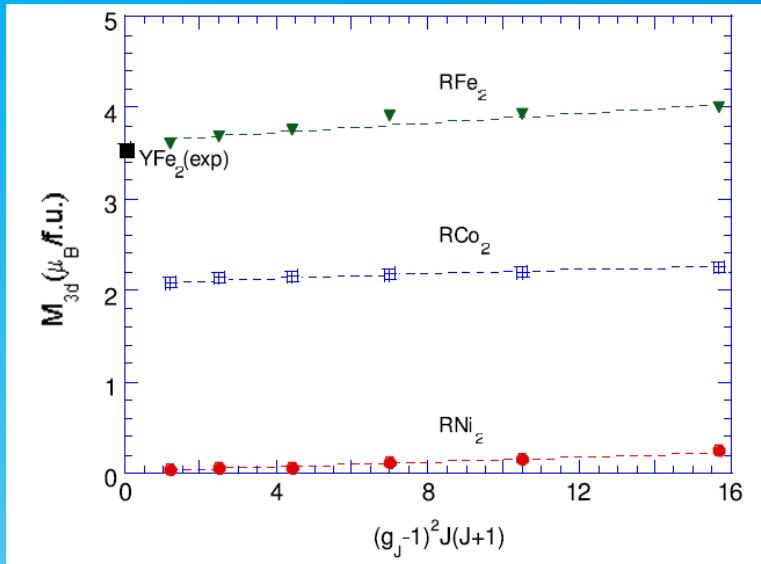
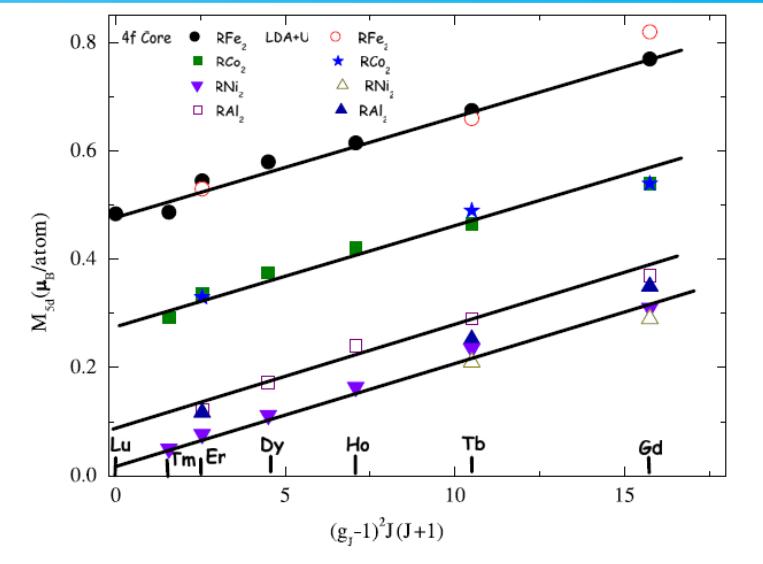
### 2.1 Exchange interactions $RCo_2$ compounds



Strong hybridization Co3d bands at site Co3b  
and Co9e with Er5d one

# 4f-5d-3d model Campbell 1972

Burzo et al, J. Phys. Cond. Matter. 23, 1026001 (2011)



$$M_{5d} = M_{5d}(0) + \alpha G \quad G = (g_J-1)^2 J(J+1)$$

4f-5d

$$J_{4f-5d} = \int g(\rho(r)) \phi_{4f}^2(r) \phi_{5d}^2 dr$$

$$M_M = M(0) + \alpha G$$

$$5d-3d \text{ short range exchange interactions}$$

$$H = -2J_{3d-5d} S_{5d}(0) \sum_i S_{3d_i}$$

↓

$$M_{5d}(0) \propto \sum z_i M_i \quad \frac{M_{5d}(0)}{\sum z_i M_i} = 2 \cdot 10^{-2}$$

$$\begin{aligned} M_{5d}(0) &= 0.03 \mu_B R\text{Ni}_2 \\ &= 0.29 \mu_B R\text{Co}_2 \\ &= 0.49 \mu_B R\text{Fe}_2 \end{aligned}$$

# Parimagnetism, Griffiths phase

$T > T_c$

The 5d-3d coupling exist at  $T > T_c$

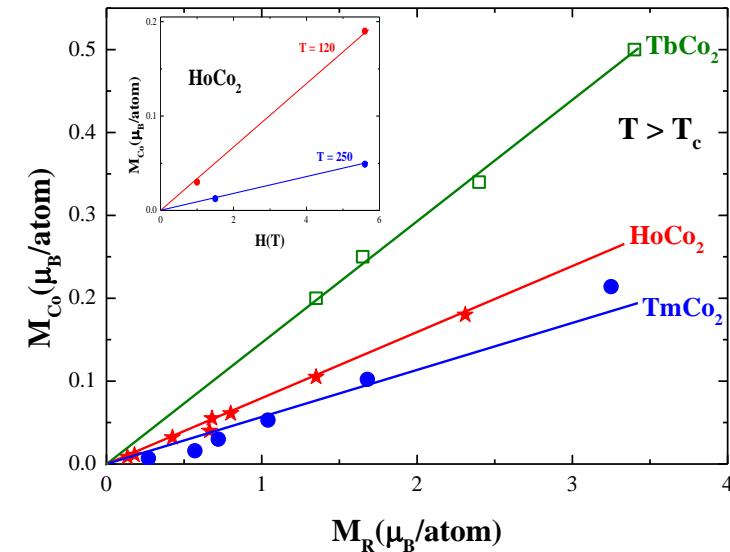
$H_{ext}$  align  $M_R$  moments



$J_{5d-3d}$  coupling

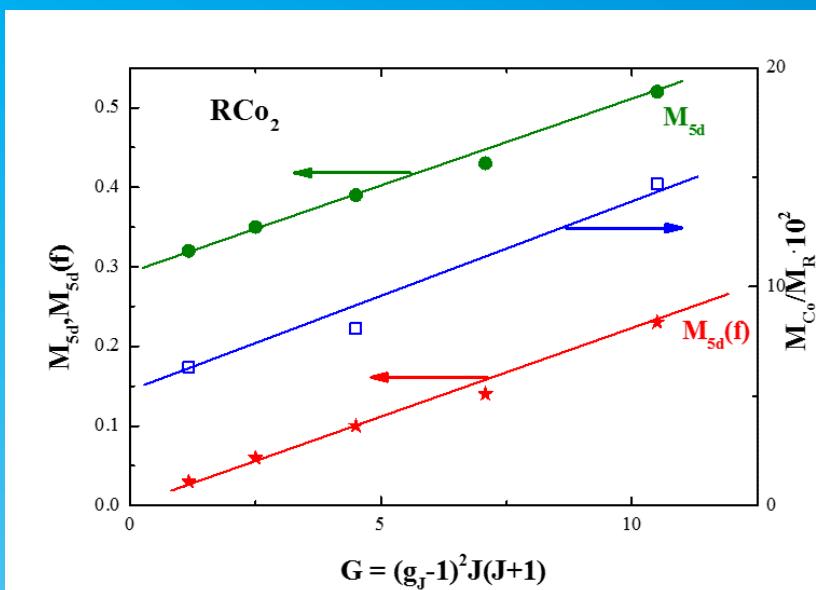
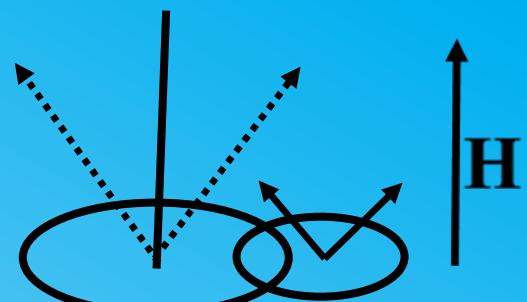
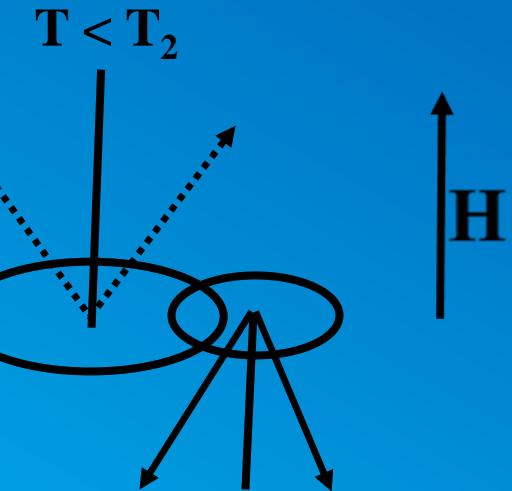
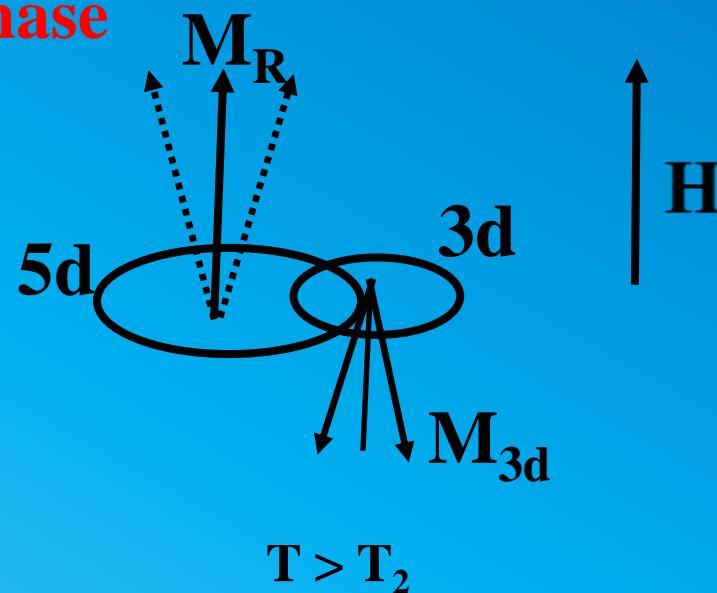


Antiparallel oriented Co moment



$$J_{4f-5d} \propto G; G = (g_J - 1)^2 J(J+1)$$

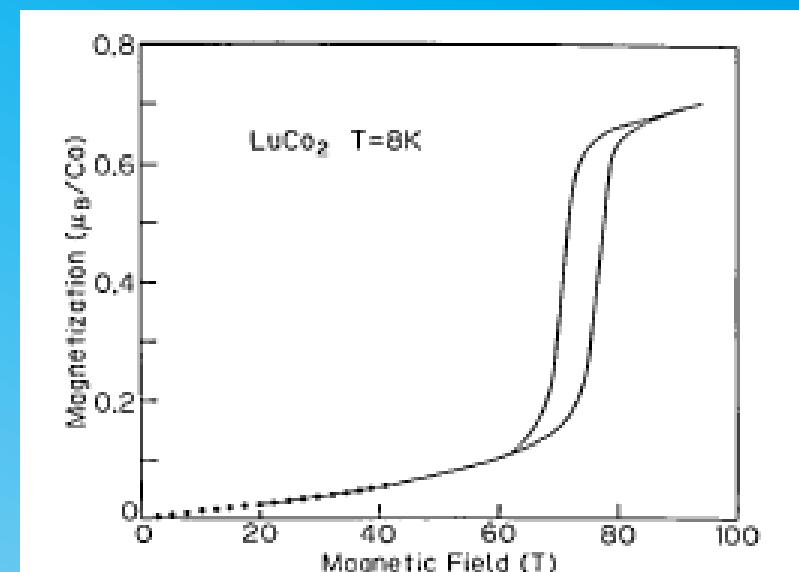
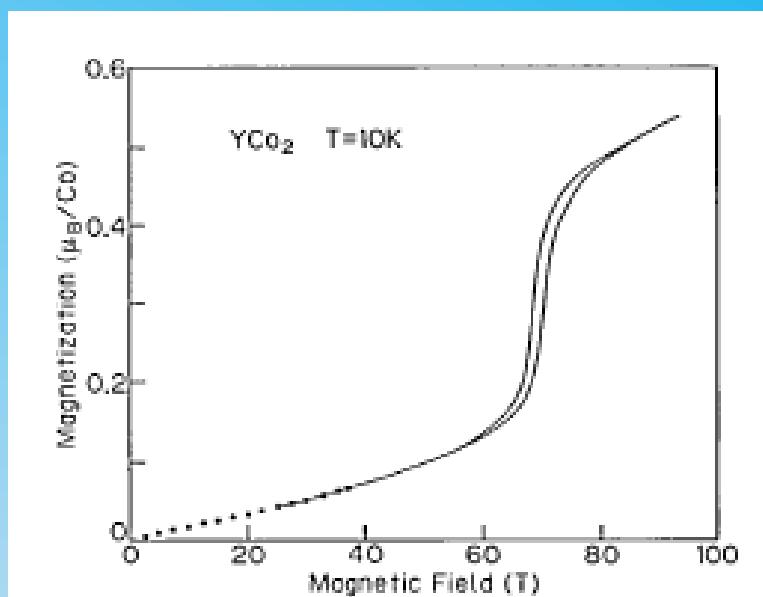
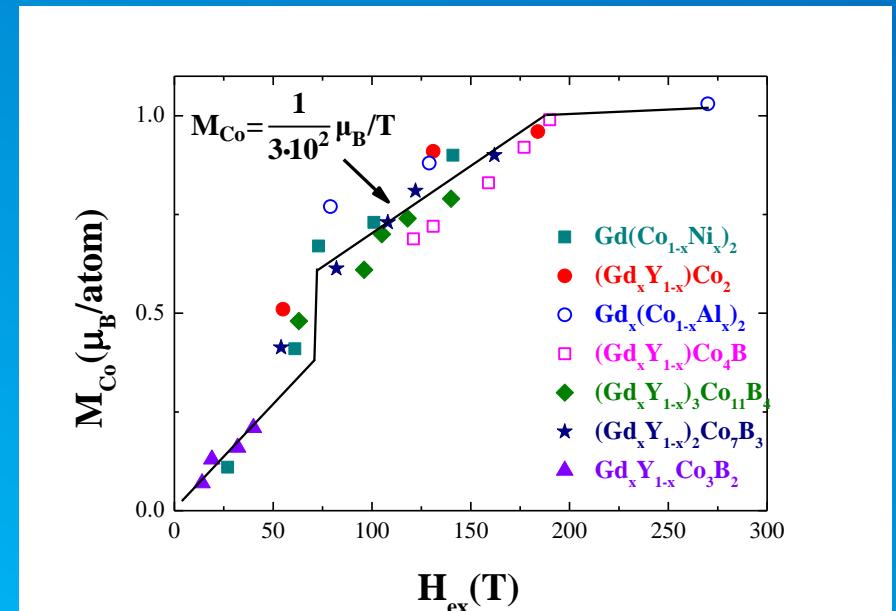
$$M_{5d}(f) = \alpha G, \quad \alpha = 2.1 \cdot 10^{-2} \mu_B$$



## 2.2 Induced cobalt moments

Critical field for inducing cobalt ordered moment

- Itinerant electron metamagnetism: conditions for a paramagnetic substance to become ferromagnetic by application and subsequent removal of strong magnetic field (Wholfarth-Rhodes, 1962).
- Induced magnetism (epamagnetism): shift of the spin-up and spin down bands under the action of exchange of external field (Burzo 1977).



# $\text{RCo}_4\text{M}$ , M = Ga, Si, Al

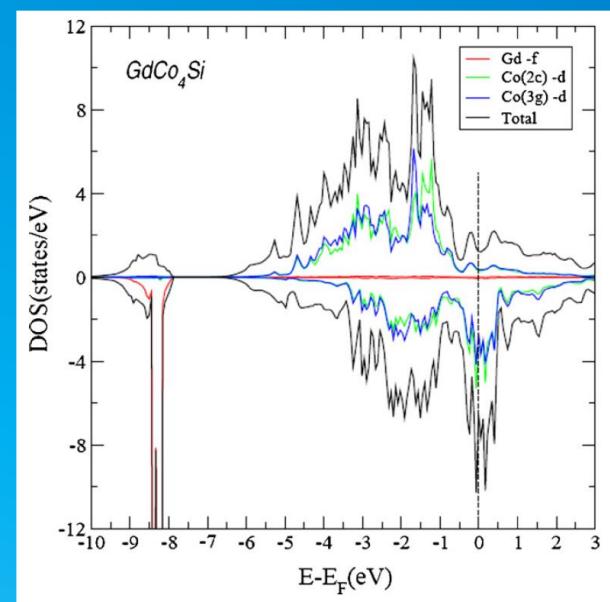
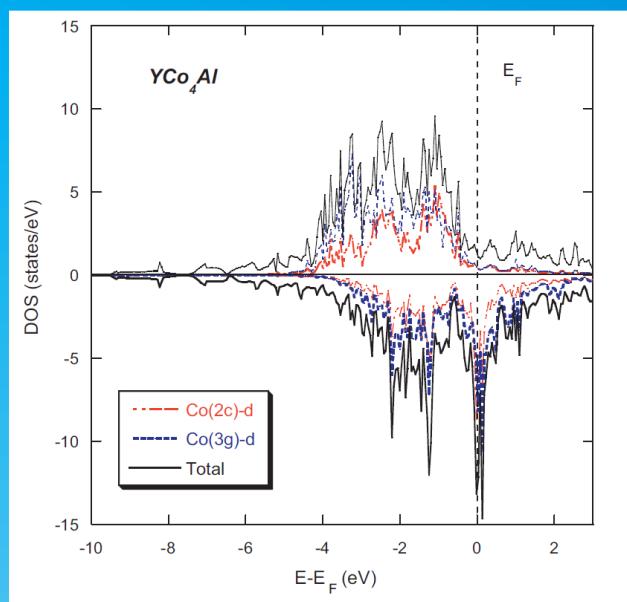
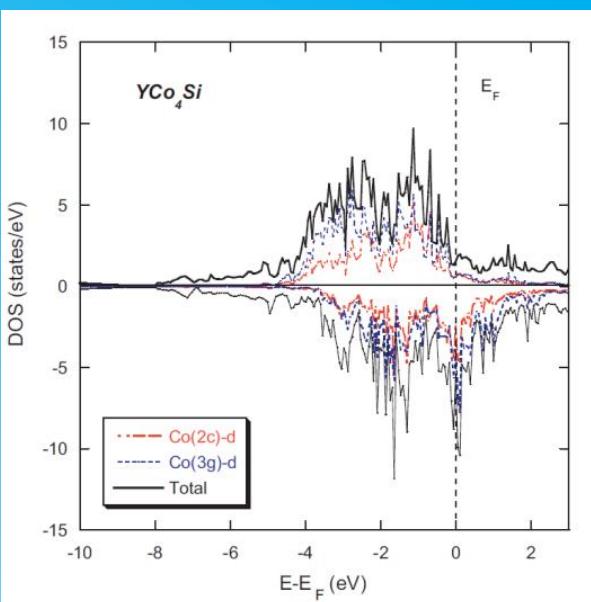
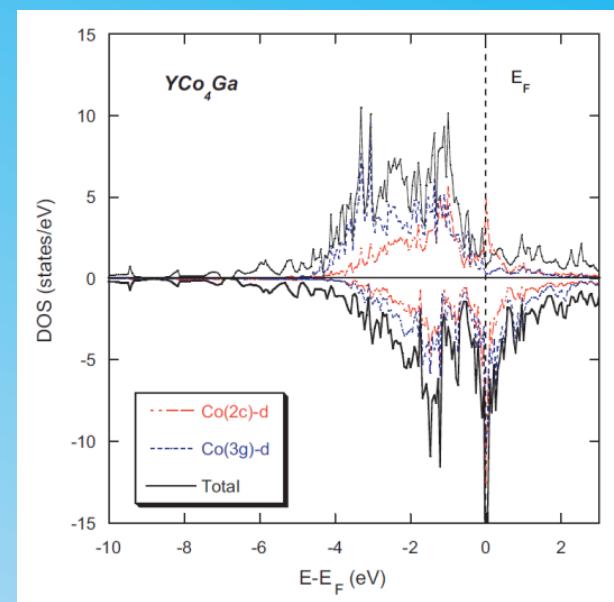
$\text{M}_{\text{Co}}$  strongly dependent on composition

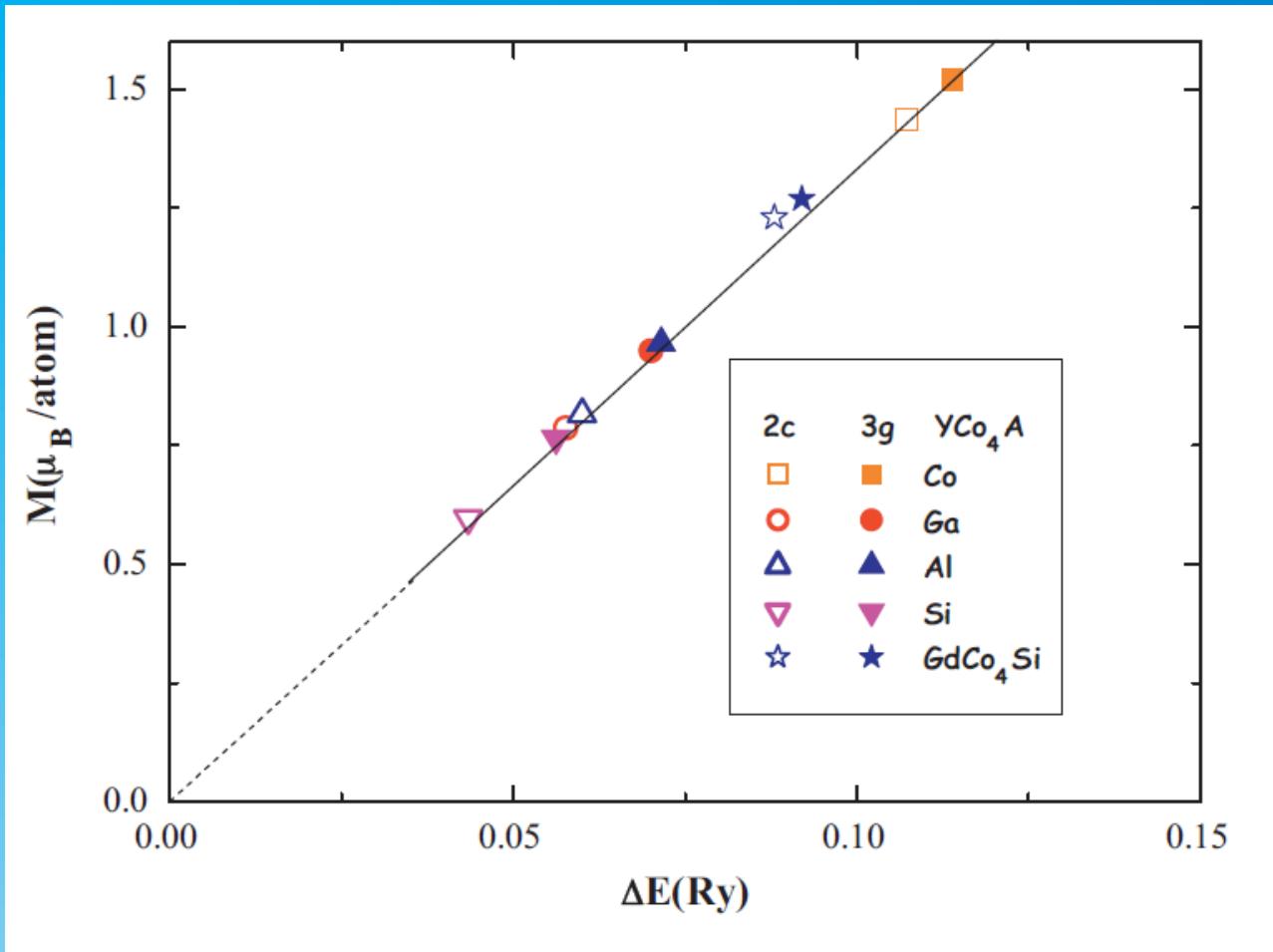


exchange interactions



splitting on 3d bands





$s = 0.95 \text{ eV}/\mu_B$  general characteristic

No presence of itinerant electron metamagnetism

## 2.3 Magnetism and pressure effects

### Magnetic behaviour of $\text{RCO}_5$ under pressure

#### $\text{RCO}_5$ compounds

Crystal structure  $\text{CaCu}_5$

Co 2c, 3g

R1a

R = Pr, Nd, Sm, Y ferromagnetic

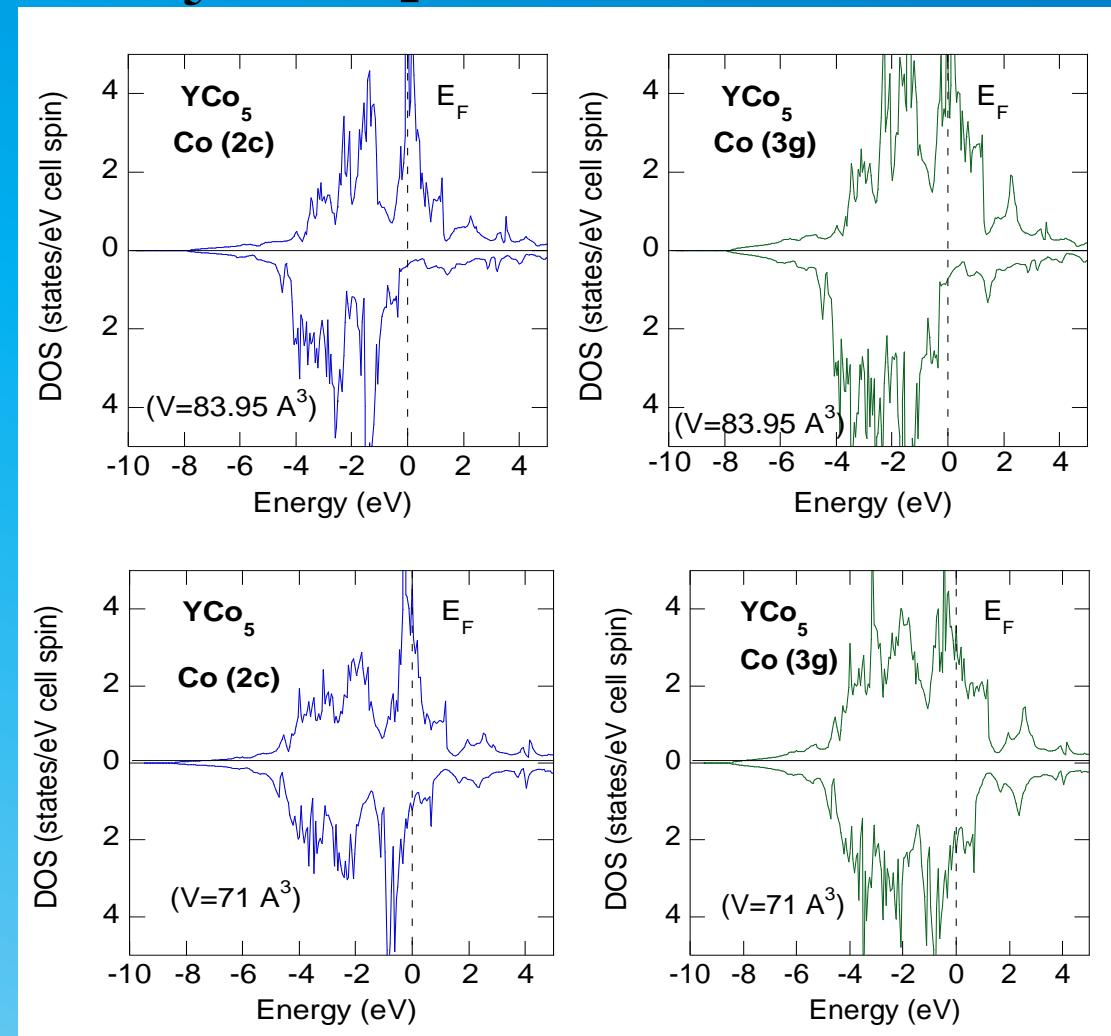
R = Gd, Tb, Dy, Ho, Er ferrimagnetic

#### $\text{YCO}_5$ , $\text{GdCO}_5$

Cobalt shows strong ferromagnetism

#### $\text{YCO}_5$

$H_a = 18 \text{ T}$  at  $4 \text{ K}$ ,  $T_c = 1000 \text{ K}$



spin-up sub-band

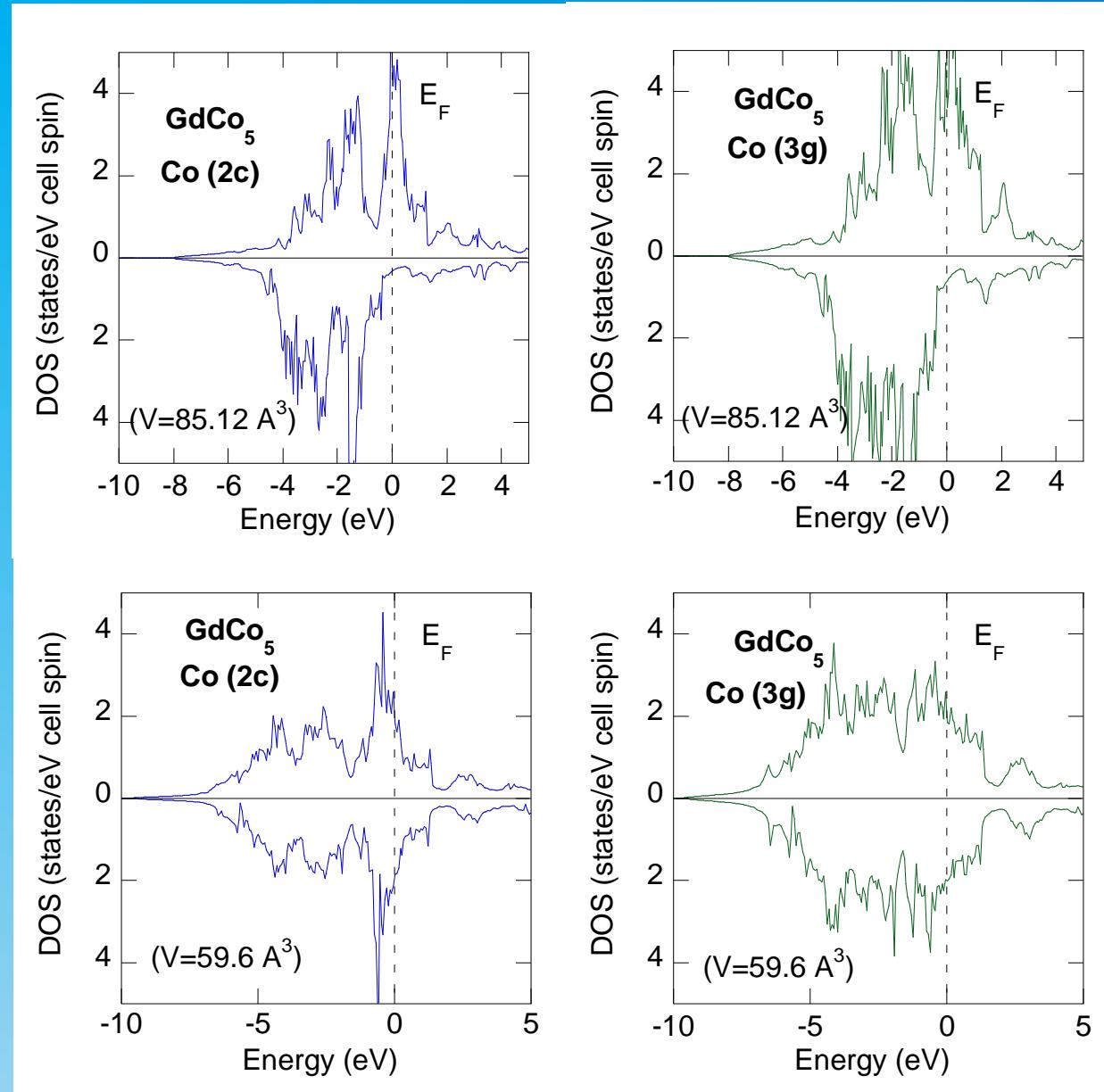
$\rightarrow$  shift to lower energies

spin-down sub-band

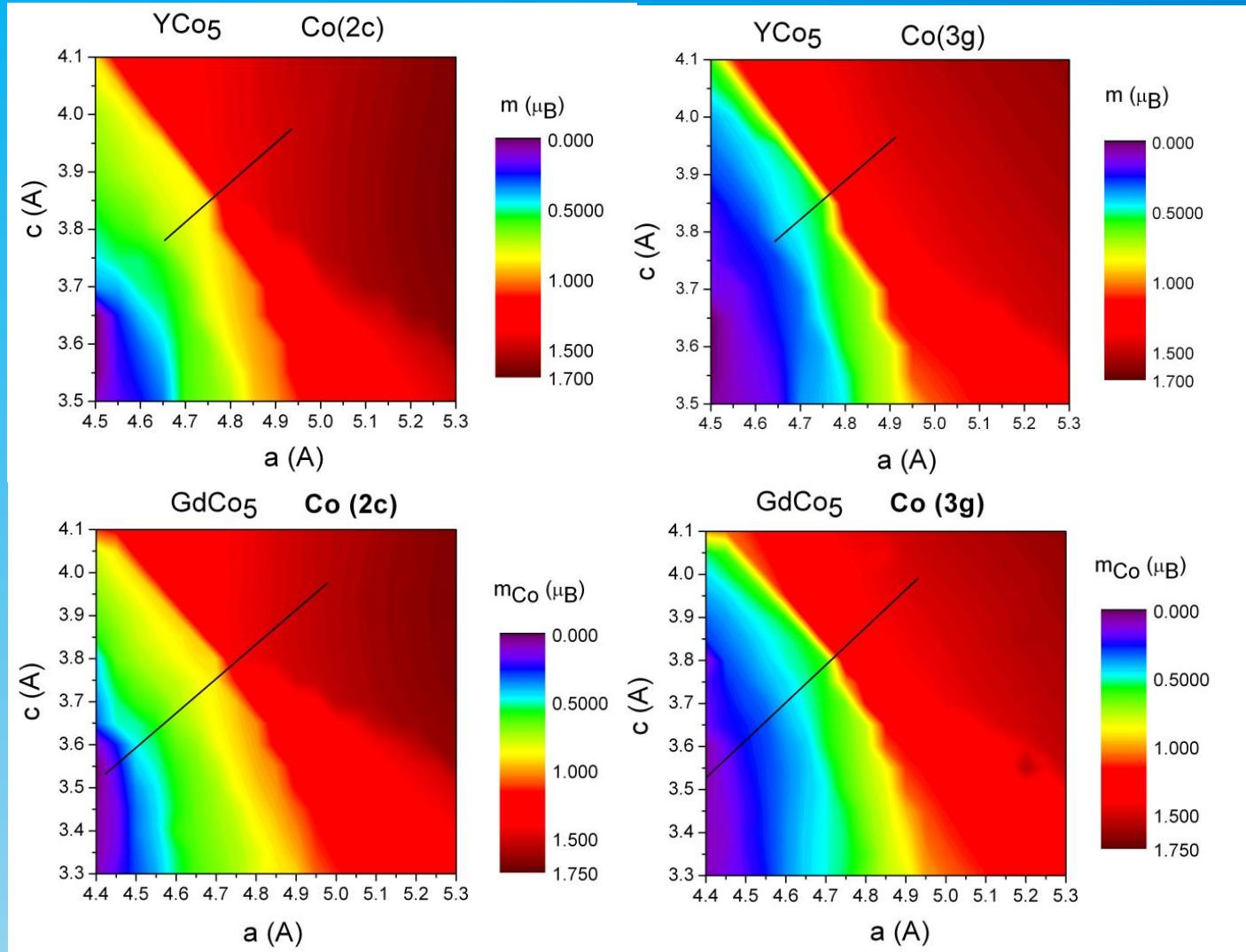
$\rightarrow$  shift to higher energies

} exchange splitting diminishes

# **GdCo<sub>5</sub>: shift of sub-bands as in YCo<sub>5</sub>**



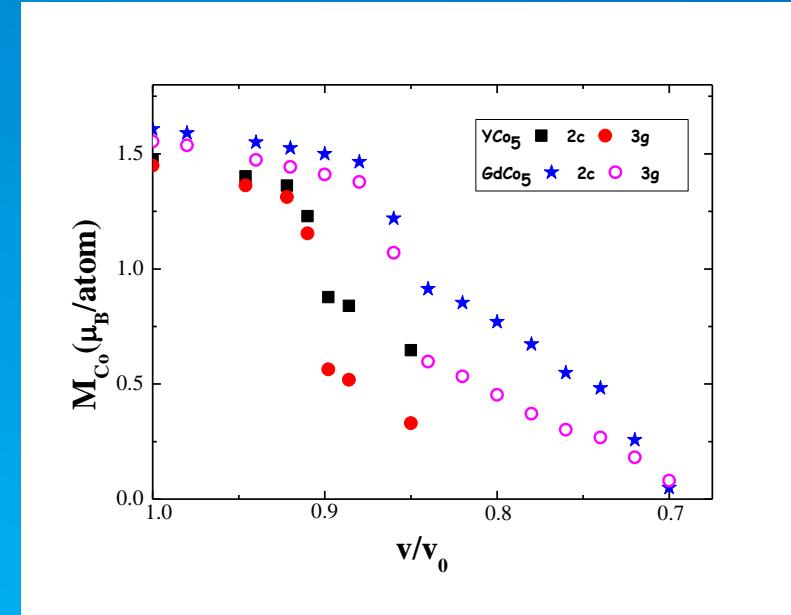
# Dependence of cobalt moment on lattice parameters



## Transition from high spin state (HS) to low spin state (LS)

$\text{YCo}_5 \quad v/v_0 = 0.90$

$\text{GdCo}_5 \quad v/v_0 = 0.86$



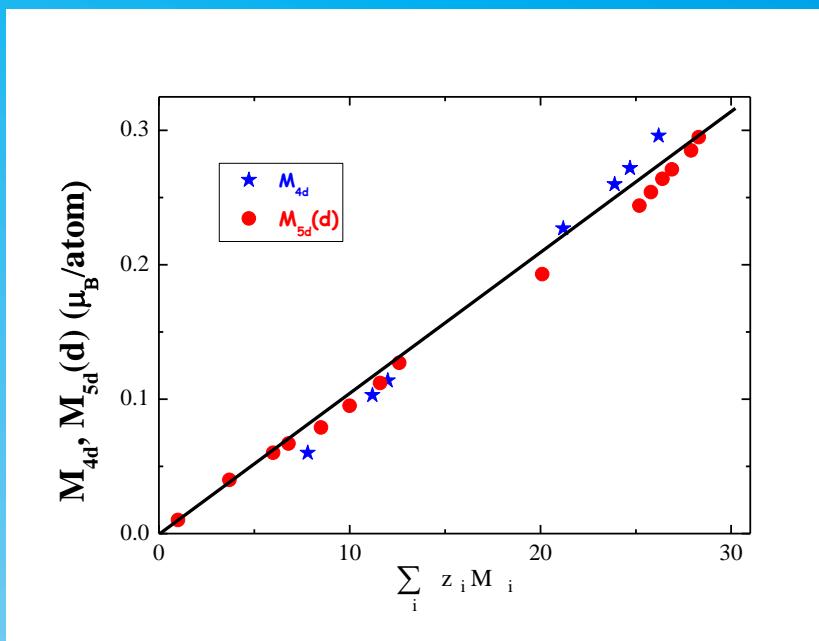
R-Co 4f-5d-3d exchange interactions

Y-Co 4d-3d exchange interactions

$$M_{5d} = M_{5d}(0) + M_{5d}(f)$$

$$M_{5d}(0) \propto \sum z_i M_{\text{Co}_i}; \quad M_{4d} \propto \sum z_i M_{\text{Co}_i}$$

The same behaviour for  $M_{5d}(0)$  and  $M_{4d}$



## **Pressure effects**

### **Co strong ferromagnet:**

- high hydrostatic pressure destroys the strong ferromagnetism
- the transformation proceeds in a stepwise fashion concomitant with isomorphic lattice charge
- there is a shift of the bands
  - spin up to lower energies
  - spin down to lower energies

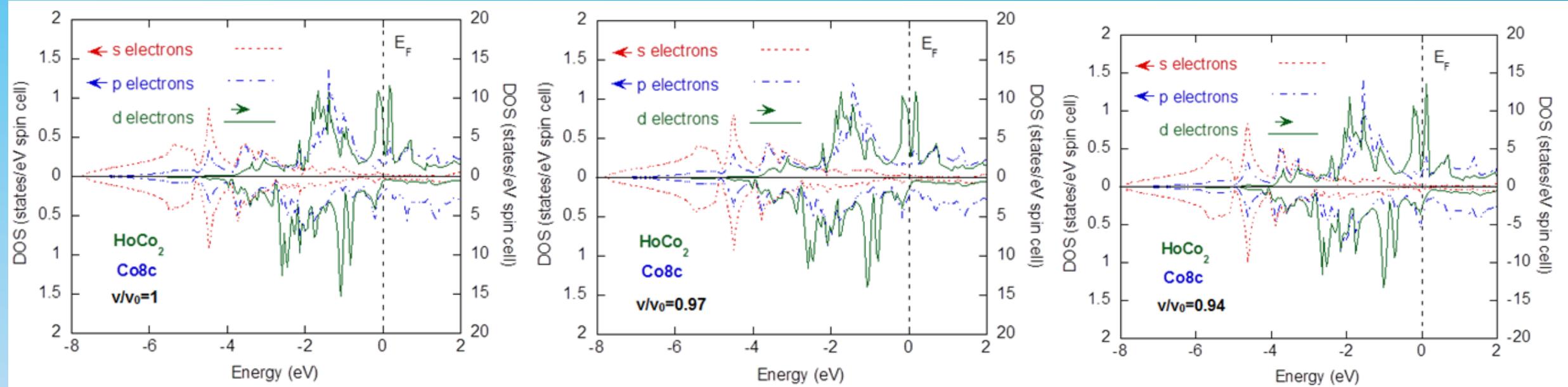
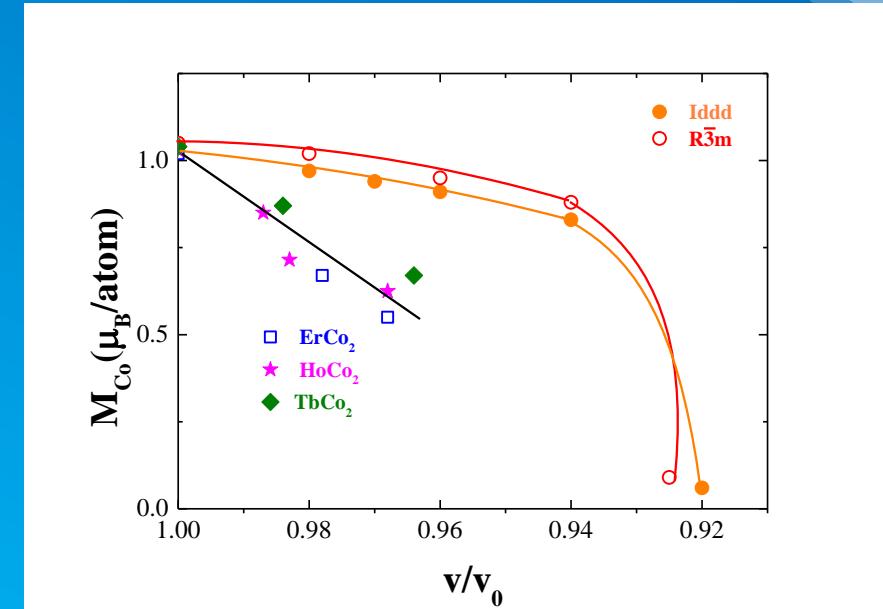
### **Unstable thermodynamic state**

rather high DOS at both spin up and spin down bands

## 2.3.2 Weak ferromagnetism Pressure effects



$T < T_c = 85 \text{ K}$  Crystal structure: tetragonal  $I4_1/\text{amd}$



# Magnetovolume effects

## $\text{RCO}_2$

$$\Gamma = \frac{1}{k_B T_c} \frac{dT_c}{dp} = \frac{d \ln T_c}{d \ln v}$$

$$\Gamma = \frac{5}{3} + BT_c^{-2} \quad \text{band model}$$

$$B = \frac{5}{6} \frac{I}{I_b} I T_F^2, \quad T_c = T_F (\bar{I} - 1)^{-1}$$

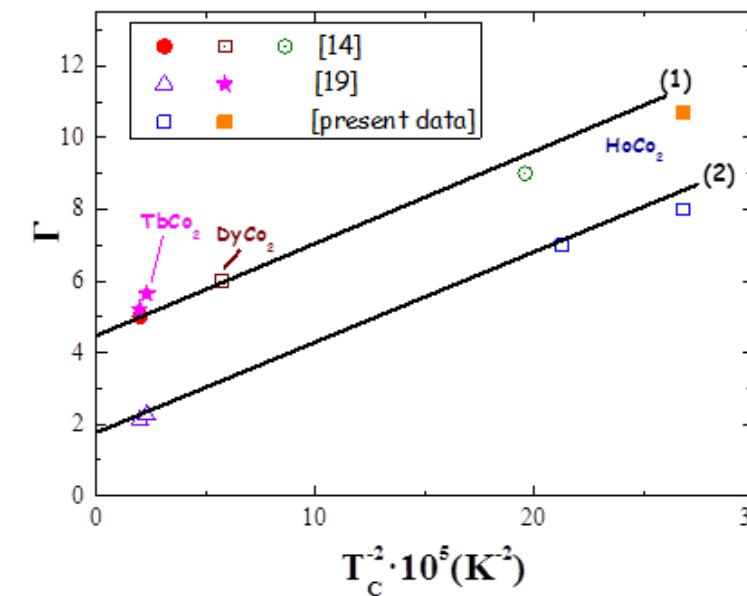
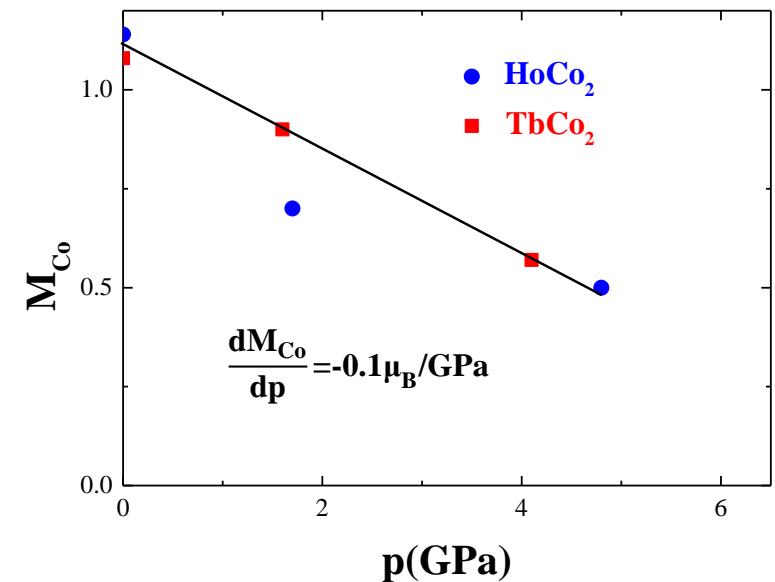
I effective intra-atomic exchange integral reduced from its bare value  $I_b$ ,  
 $\bar{I} = I \eta(E_F)$

$$B = 3.5 \cdot 10^4 T_c^{-2}, \quad T_F = 240 \text{ K}$$

$$I/I_b = 0.85$$



important correlation between 3d electrons



# $\text{RCO}_2$ compounds

Behaviour: spin fluctuations

continuous change of the degree of localization of the cobalt moment

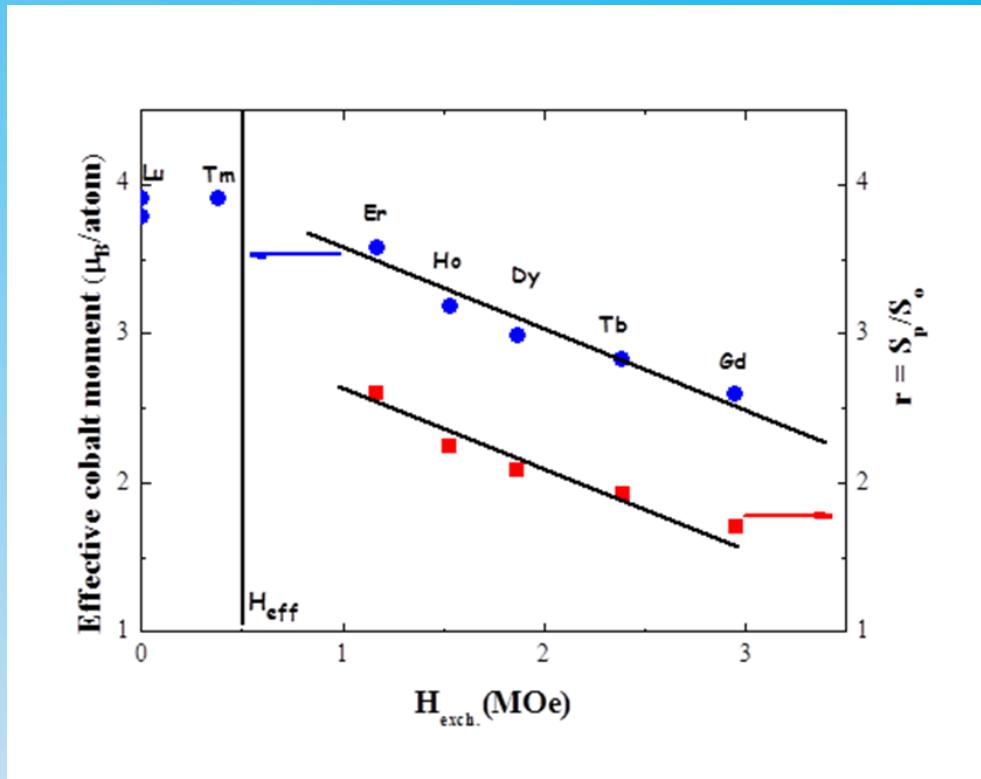
$$r = S_p/S_0$$

$$\mu_{\text{eff}} = g \sqrt{S_p(S_p+1)}$$

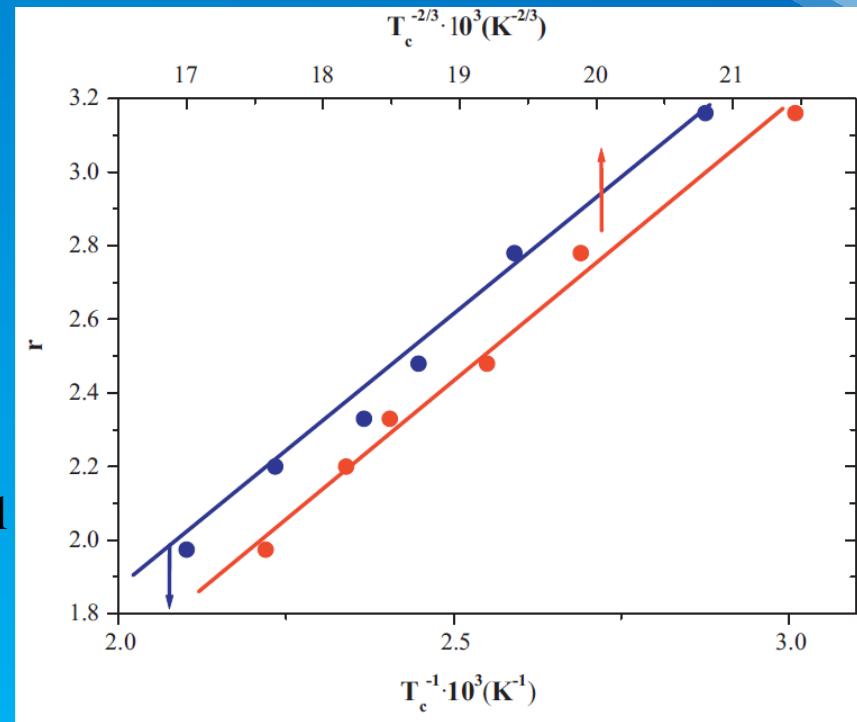
$$\mu_0 = gS_0$$

$r \propto T_c^{-2/3}$  spin fluctuations

$r \propto T_c^{-1}$  band model



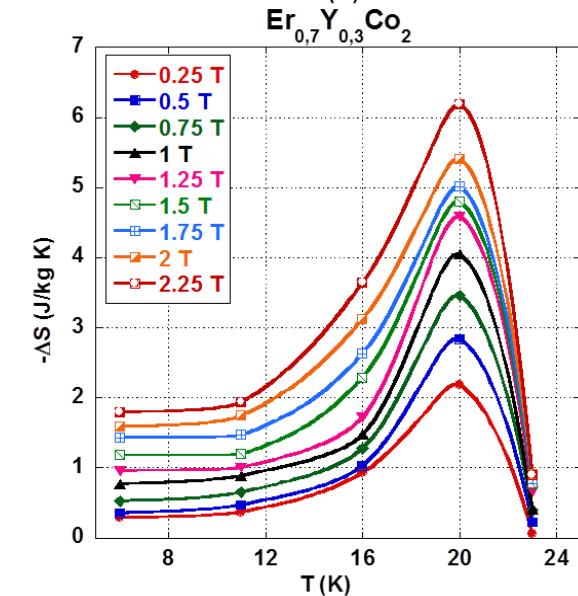
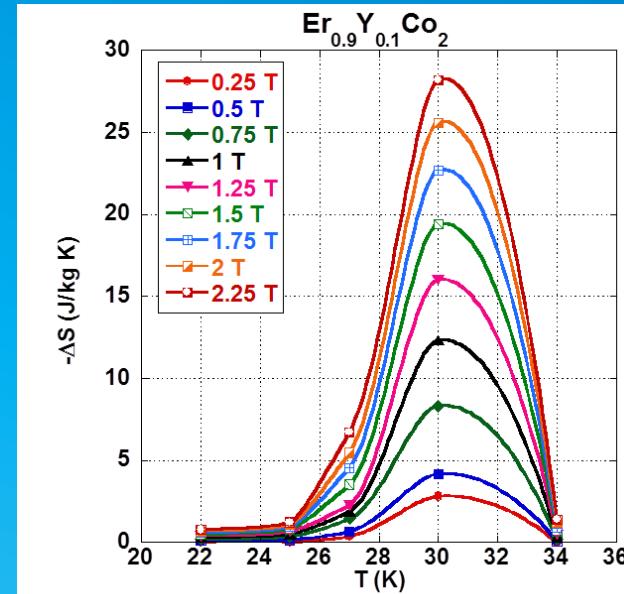
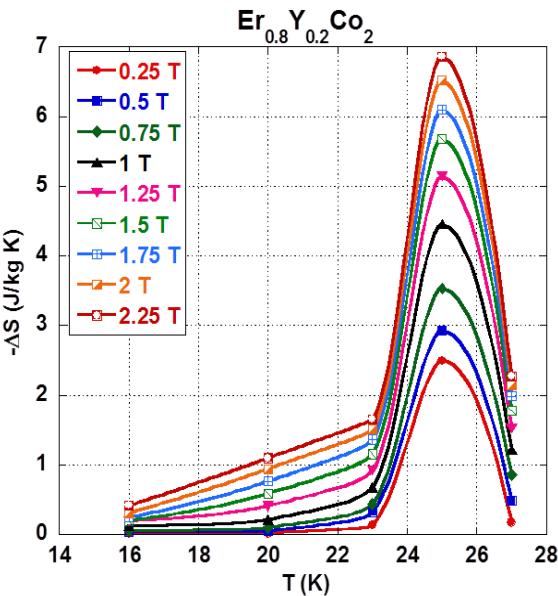
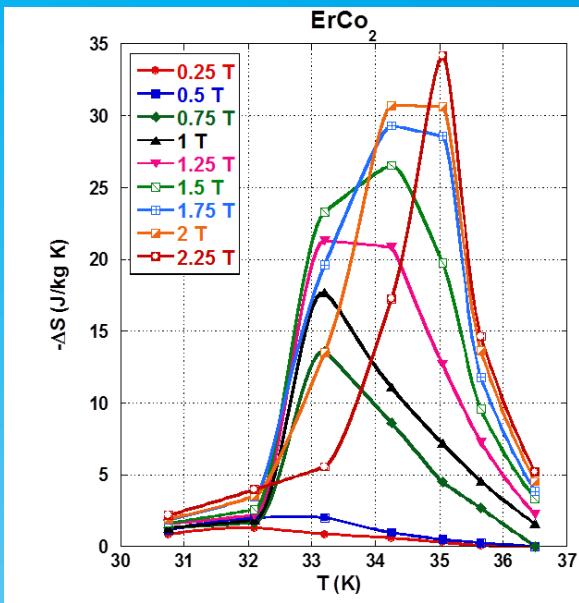
$r = 1$  localized model  
 $r \rightarrow \infty$  band model



The degree of localization increase with  $H_{\text{exch}}$

# Magnetocaloric effect

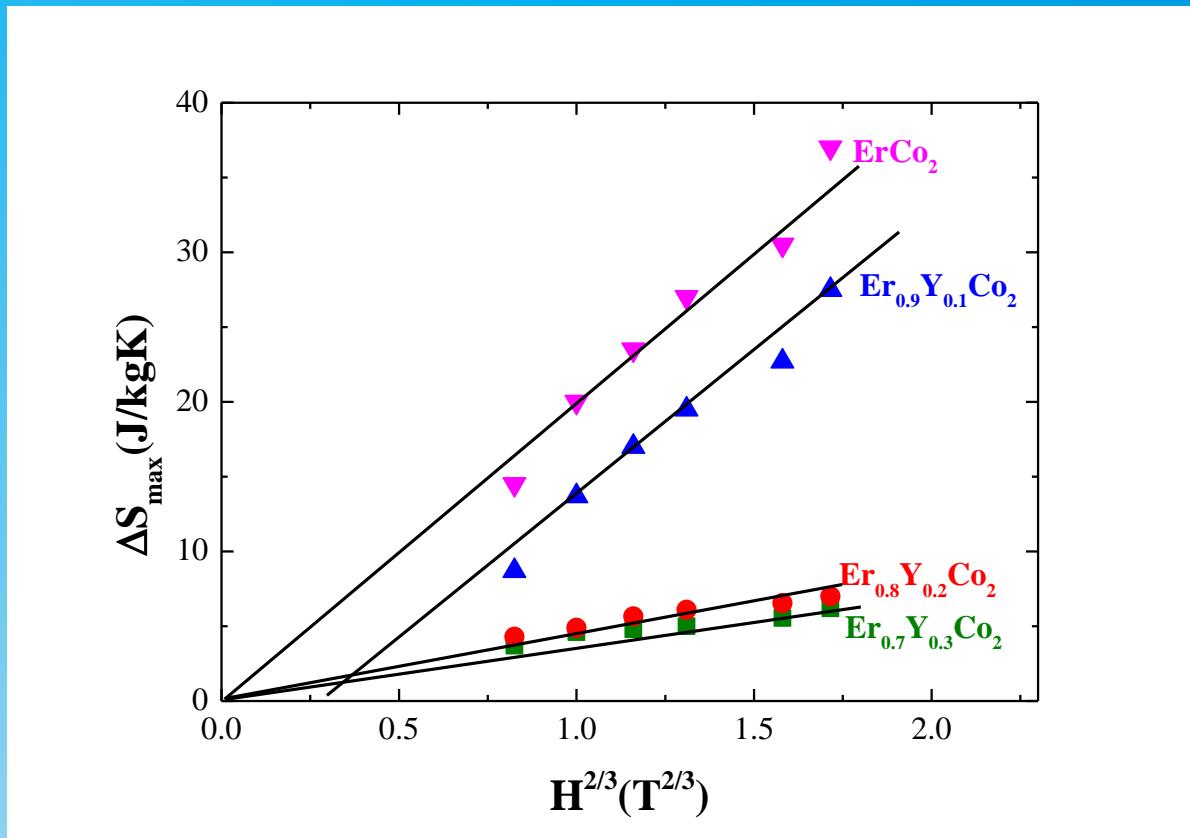
$\text{Er}_{1-x}\text{Y}_x\text{Co}_2$



# $\text{Er}_{1-x}\text{Y}_x\text{Co}_2$

High magnetocaloric effect: first order transition  $x = 0, x = 0.1$

Low magnetocaloric effect: second order transition  $x = 0.2, x = 0.3$

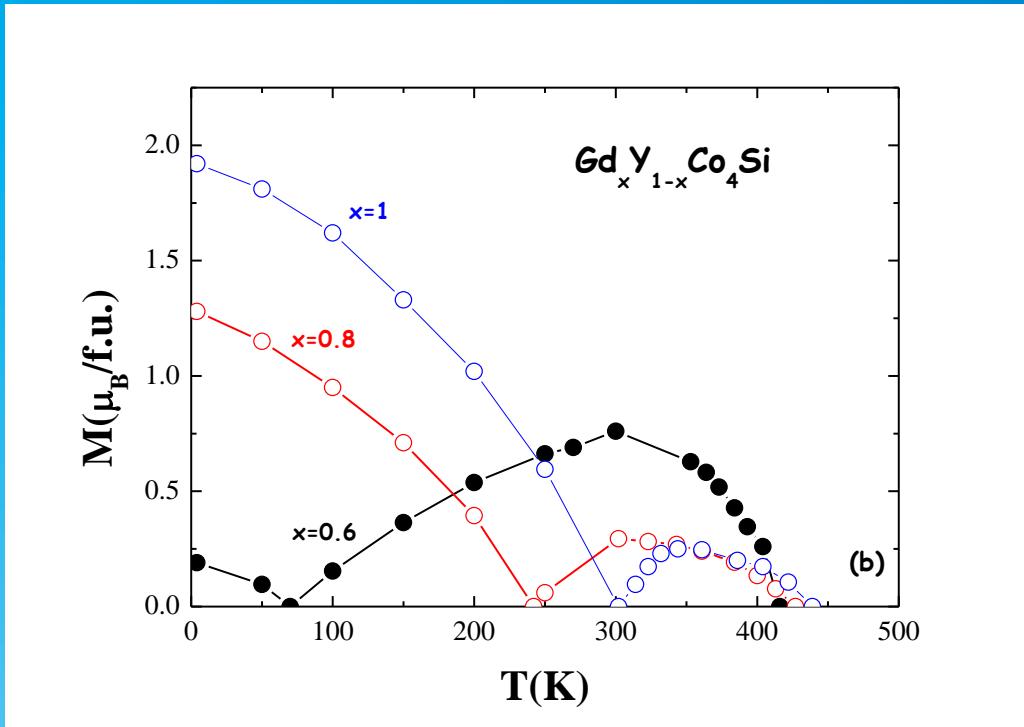
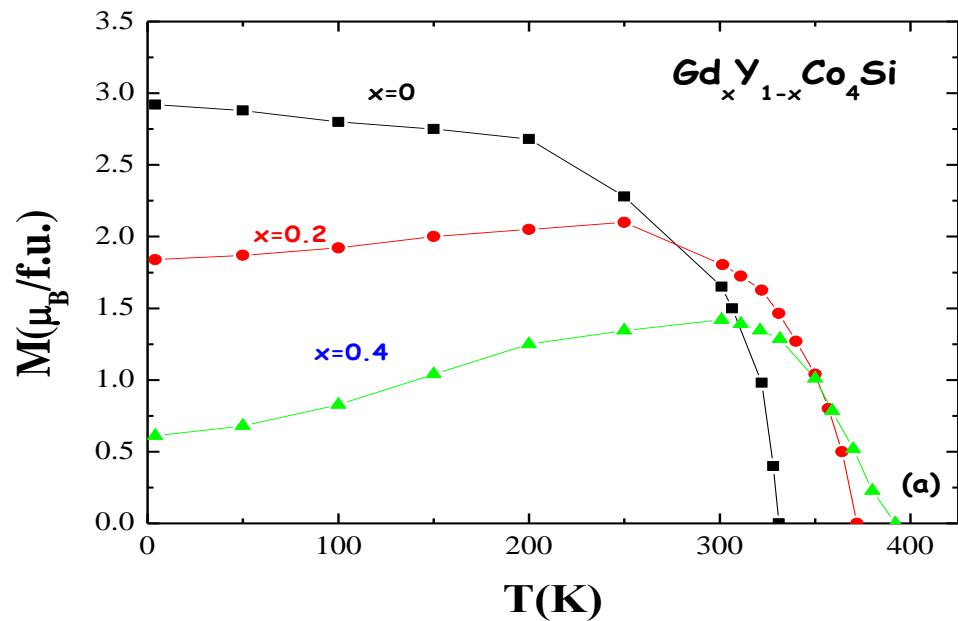


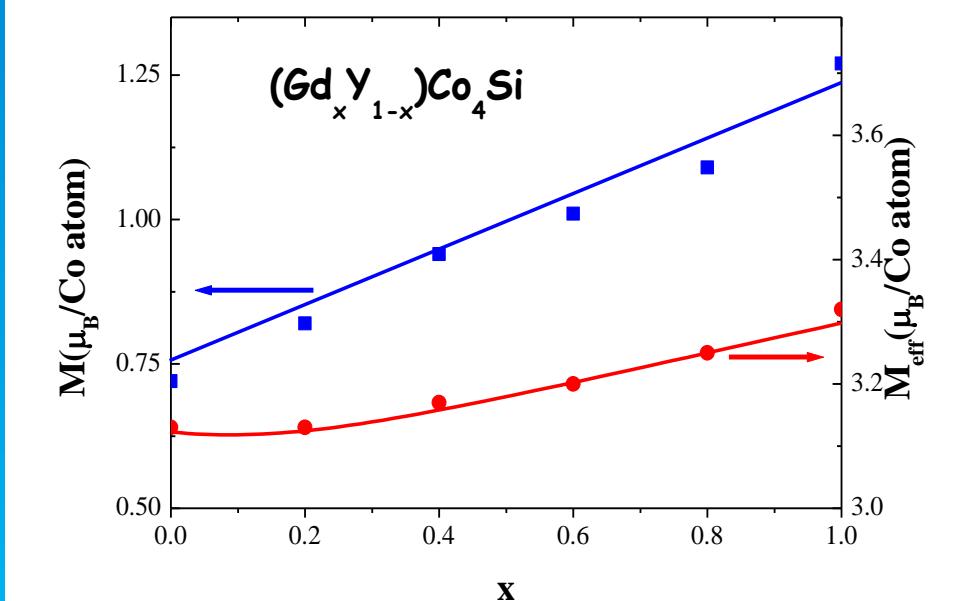
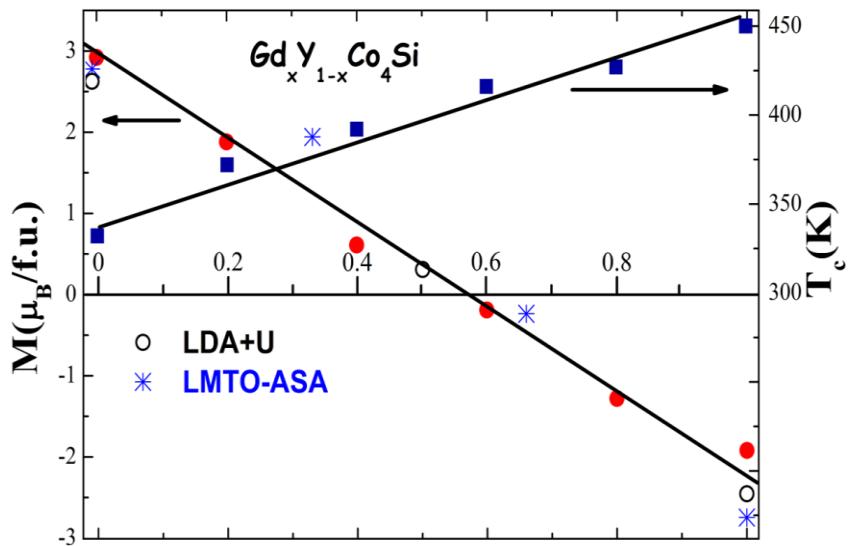
$$\Delta S \propto H^{2/3}$$

## 2.4 Weak to strong ferromagnetism

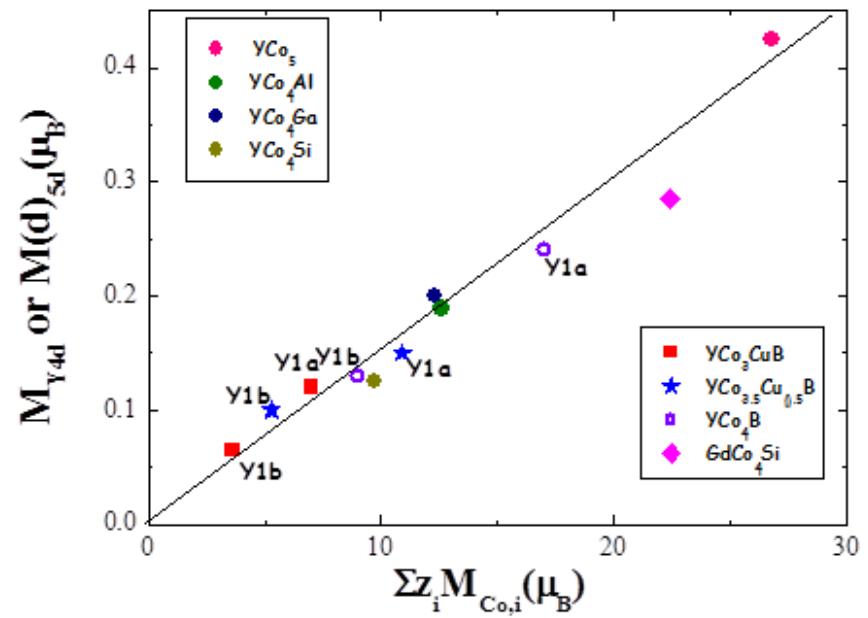
### $Gd_x Y_{1-x} Co_4 Si$ , Gd-Co-B

Transition from strong ferromagnetism to weak ferromagnetism. Example  $Gd_x Y_{1-x} Co_4 Si$ .





$$M_{5d}(0) \text{ or } M_{5d} \propto \sum_i z_i M_{Co_i}$$



# $\text{RCO}_4\text{B}$ , R = Ce, Y, Gd, $\text{R}_2\text{Fe}_{17}\text{C}_2$

## Pressure effects

Localized model

$$a = -\frac{5}{3} + \frac{\text{dln}J_{\text{eff}}}{\text{dln}v}$$

$$b = \frac{5}{8} \frac{k_B N_0 g^2 I^2}{S(S+1) J_{\text{eff}}^2 I_b}$$

I effective intra-atomic exchange integral reduced from bare value  $I_b$

$$\text{RCO}_4\text{B} \quad a = 13.4 \quad b = -0.022 \text{ K}^{-1}$$

$$\text{R}_2\text{Fe}_{17}\text{C}_2 \quad a = 37.5 \quad b = -0.0063 \text{ K}^{-1}$$

$$\frac{\text{dln}J_{\text{eff}}}{\text{dln}v} \begin{cases} 6 \\ 11 \end{cases} \quad \begin{matrix} \text{RCO}_4\text{B} \\ \text{R}_2\text{Fe}_{17}\text{C}_2 \end{matrix}$$

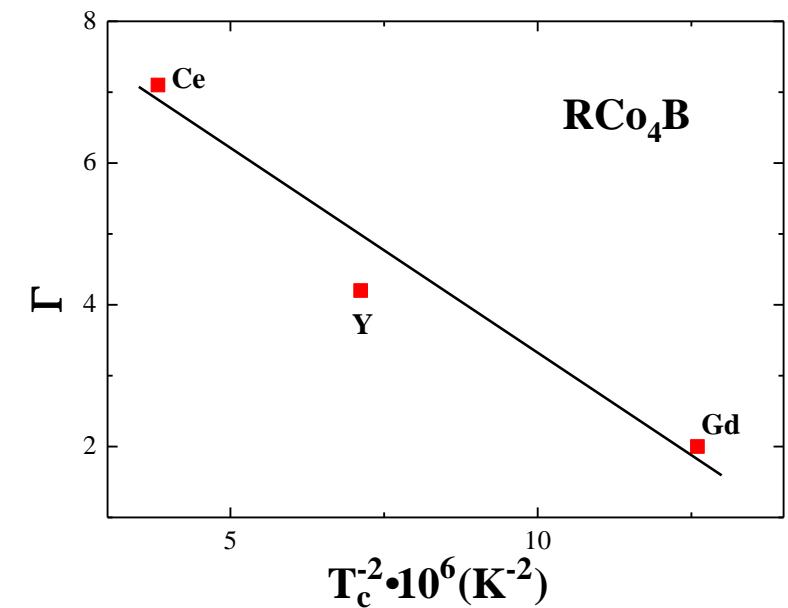
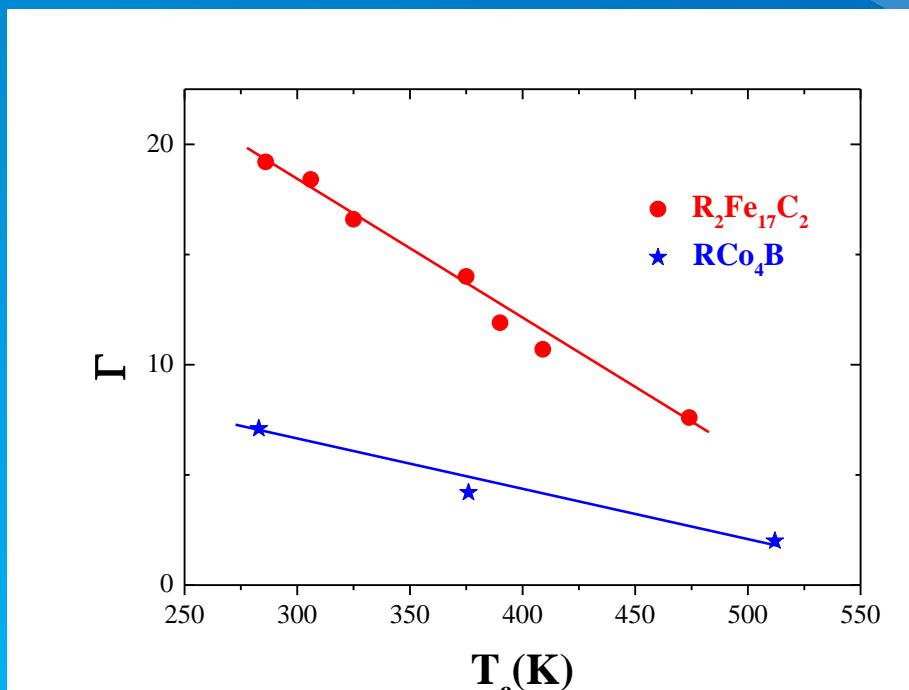
$\text{RCO}_4\text{B}$  can also be described by

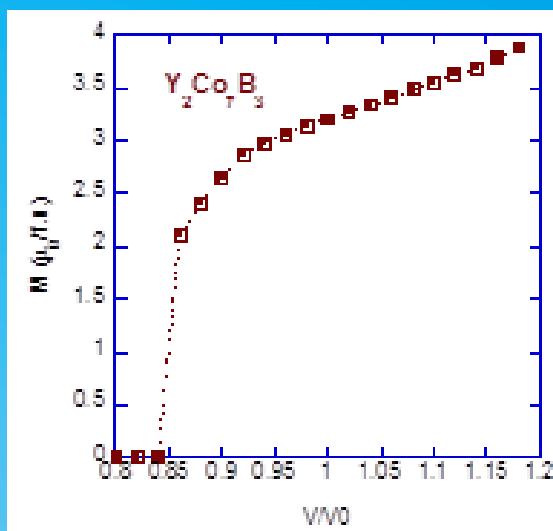
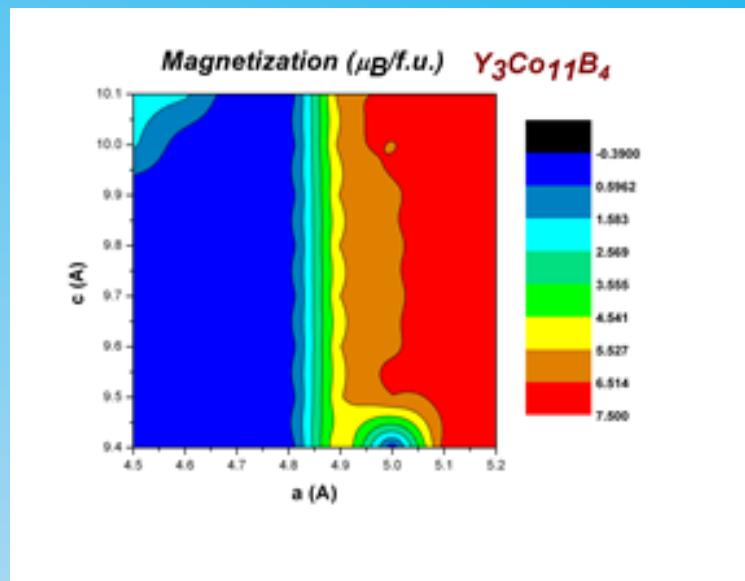
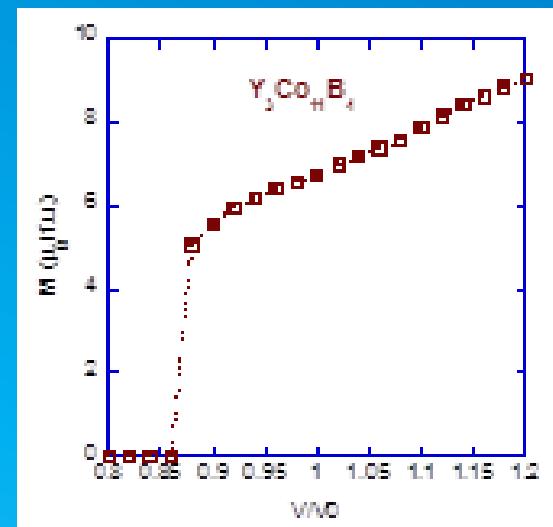
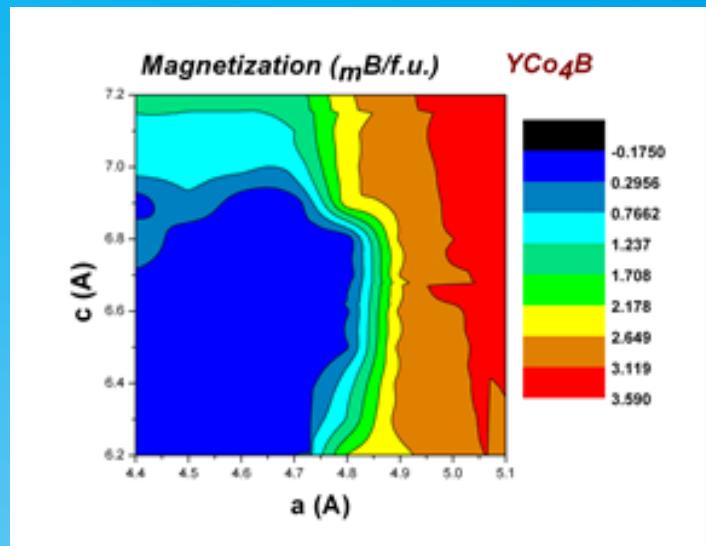
$$\Gamma = a_1 + B T_c^{-2}$$

Weak  $\rightarrow$  strong ferromagnetism

SF model

$$\Gamma = a - b T_c$$





## *Conclusions*

Exchange enhanced paramagnets  
induced temperature moments

$$\chi \propto T^2 \rightarrow \chi \propto T^{-1}$$

Exchange interactions 4f-5d-3d model

$$M_{5d} = M_{5d}(0) + M_{5d}(f)$$

$$M_{5d}(0) \propto \sum z_i M_i \quad M_{5d}(f) \propto (g_J - 1)^2 J(J+1)$$

Induced cobalt moment

$$H_{ex} = H_{cr} \approx 70 \text{ T}$$

$$H > H_{cr} \quad M_{co} = a H_{exch} \quad a = (3 \cdot 10^{-2}) \mu_B / T$$

At  $T > T_c$ : induced moments by alignment of  $M_R$  by external field due to 5d-3d coupling.

High magnetocaloric effect: compounds showing first order magnetic transition

Pressure effect:

weak ferromagnet: direct collapse of  $M_{Co}$

strong ferromagnet: sequential collapse of  $M_{Co}$  with at least one step

## **ACKNOWLEDGMENTS**

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*Thank you very much for  
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